

Equations of Motion of Spinning Particles (in Gravity)

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- **Spinning Particles are non-geodesic, violate EEP**
- **New Action Principle for Classical Spin where $\delta\dot{x} \neq d \delta x/dt$**
- **Spin contributions to source of gravity**

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<http://www.clifford.org/~wpezzag/talks.html>

I. Outline

- Title/Introduction 1
- **I. Outline** 2

- **II. Hierarchal Handwaving**
 - A. Spin Invariant Mass 3
 - B. Use Geometry to Unify Equations 4
 - C. Polydimensionalism 5

- **III. Classical Mechanics of Spin**
 - A. New Affine Parameter 6
 - B. Papapetrou Equations as Autoparallels 7
 - C. Lagrangians in Curved Space 8

- **IV. Anholonomic Mechanics**
 - A. The Secret Ingredient: Vel of Var is not Var of Vel 9
 - B. Hamiltonian Mechanics 10
 - C. Spin as source of gravity 11

- **V. References** 12

II.A Hierarchal

Handwaving

3 ⊕ 1 Special Relativity

$$(mc)^2 = E^2/c^2 - \vec{p}^2$$

invariant rest mass (Scalar)² - (Vector)²

Motion increases Mass

$$m' = m \sqrt{1 + \left(\frac{\vec{p}}{mc}\right)^2}$$

Approx: $m' \approx m + \frac{1}{2} \frac{\vec{p}^2}{mc^2}$

4-vector $p = p^u e_u$

invariant $(mc)^2 = \|p\|^2 = p_u p^u$

Classical Spin Mechanics

Dixon (70) Mass is Dynamic

$$(m_0 c)^2 = p_u p^u - \frac{S^{uv} S_{uv}}{2\lambda^2}$$

invariant mass (Vector)² - (Bivector)²

Spin increases Mass

$$m = m_0 \sqrt{1 + \left(\frac{|S|}{m_0 c \lambda}\right)^2}$$

Approx: $m \approx m_0 + \frac{1}{2} \frac{|S|^2}{m_0 \lambda^2 c^2}$

$\lambda \approx$ Radius of gyration

Polyvector: $\tilde{m} = p^u e_u + \frac{S^{uv}}{2\lambda} e_u e_v$

invariant:

$$|\tilde{m}|^2 = p_u p^u - \frac{S^{uv} S_{uv}}{2\lambda^2}$$

II.B. USE Geometry to Unify Equations

1. Unify 3D Electrodynamics with 4D

$$\left. \begin{array}{l} \text{- Scalar: } \dot{\mathcal{E}} = e \vec{E} \cdot \vec{v} \\ \text{- Vector: } \dot{\vec{p}} = e(\vec{E} + \vec{v} \times \vec{B}) \end{array} \right\} \dot{p}^\mu = \frac{e}{m} p_\nu F^{\mu\nu}$$

2. Coordinate Independent Form

$$p = p^\mu \hat{e}_\mu$$

$$F = \frac{1}{2} F^{\mu\nu} \hat{e}_\mu \wedge \hat{e}_\nu$$

$$\dot{p} = \frac{e}{2m} [p, F]$$

where $[e_\mu, e_\nu] = 2e_{\mu\nu}$, $[e_\alpha, e_{\mu\nu}] = 2e_\alpha \cdot (e_\mu \wedge e_\nu)$

3. Unify Spin and Momentum Equations

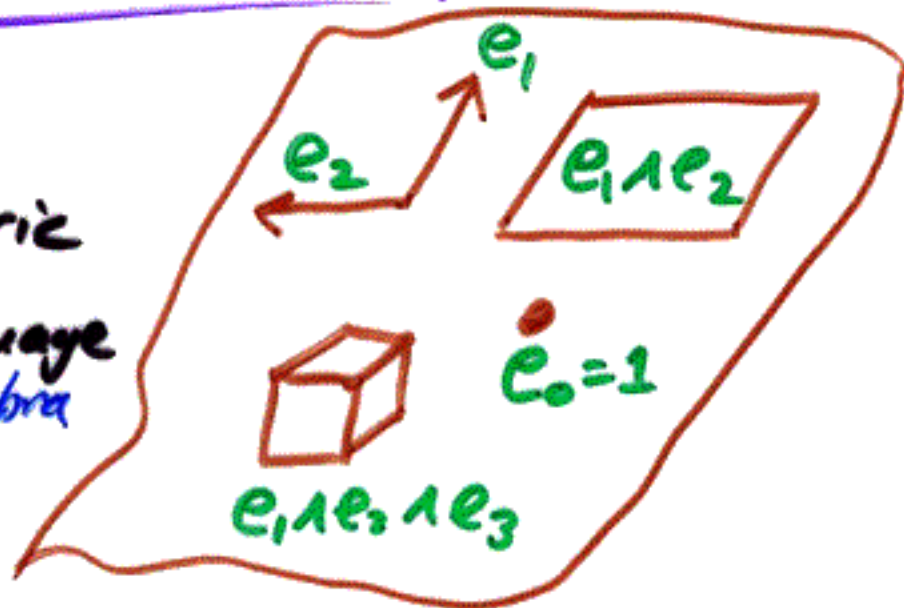
$$\left. \begin{array}{l} \dot{p}^\mu = \frac{e}{m} p_\nu F^{\mu\nu} \\ \dot{S}^{\mu\nu} = \frac{e}{m} (F^\mu{}_\lambda S^{\lambda\nu} - F^\nu{}_\lambda S^{\lambda\mu}) \end{array} \right\}$$

$$\dot{q}m = \frac{e}{2m} [q\mathfrak{m}, F]$$

Polymomenta: $q\mathfrak{m} = p^\mu e_\mu + \frac{1}{2\lambda^2} S^{\mu\nu} e_\mu \wedge e_\nu$

II.C Polydimensionalism

- ① Space is
Polygeometric
Natural Language
Clifford Algebra



② Dimensional Democracy

Each geometric element has a Coordinate

$$d\Sigma = d\phi \hat{e}_0 + dx^\mu \hat{e}_\mu + \frac{1}{2} da^{\mu\nu} \hat{e}_\mu \wedge \hat{e}_\nu + \dots$$

Event scalar vector bivector

③ Conserved Quantities Associated:

Spin : $S^{\mu\nu} = m \dot{a}^{\mu\nu}$

Momentum : $p^\mu = m \dot{x}^\mu$

Mass : $m \approx \dot{\phi}$

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hep-th/0011216

6 III.A. New Affine Parameter for Mechanics

① Define $dK^2 = dx^\mu dx_\mu - \frac{1}{2\lambda^2} da^{\mu\nu} da_{\mu\nu}$
(modulus)² = (length)² - (area)²

Total Derivative $\dot{Q} \equiv \frac{dQ}{dK} = \left(\frac{dT}{dK}\right) \frac{dQ}{dT} = \left(\frac{dT}{dK}\right) \dot{Q}$

Spin Correction Factor $\frac{dT}{dK} = \sqrt{1 - \frac{|\dot{a}|^2}{c^2 \lambda^2}} = \sqrt{1 + \frac{|\dot{a}|^2}{c^2 \lambda^2}}$

② New Action Principle: Particles follow paths for which the length swept out by the momentum minus area swept out by spin is extremum

$$\mathcal{L} = m_0 c \sqrt{\dot{x}^\mu \dot{x}_\mu - \frac{1}{2\lambda^2} \dot{a}^{\mu\nu} \dot{a}_{\mu\nu}}$$

③ Canonical Momenta has Spin Corrected MASS

$$p_\mu = \frac{\delta \mathcal{L}}{\delta \dot{x}^\mu} = m_0 \dot{x}_\mu = m \dot{x}_\mu$$

$$S_{\mu\nu} = \lambda^2 \frac{\delta \mathcal{L}}{\delta \dot{a}^{\mu\nu}} = m_0 \dot{a}_{\mu\nu} = m \dot{a}_{\mu\nu}$$

$$m = m_0 \frac{dT}{dK} = m_0 \sqrt{1 + \frac{|S|^2}{(m_0 c \lambda)^2}}$$

III.B Autoparallel Paths

① Parallel transport of basis vector

$$\frac{de_\mu}{d\kappa} = \left(\dot{\chi}^\sigma \frac{\partial}{\partial \chi^\sigma} + \frac{1}{2} \dot{\alpha}^{\alpha\beta} \frac{\partial}{\partial \alpha^{\alpha\beta}} \right) e_\mu$$

where $\frac{\partial}{\partial \chi^\sigma} (p^\nu e_\nu) = e_\nu (\nabla_\mu p^\nu) = e_\nu (\Gamma_{\mu\sigma}^\nu p^\sigma + d_{\mu\nu} p^\nu)$

Conjecture: Bivector Derivative is like a loop -

$$\frac{\partial}{\partial \alpha^{\alpha\beta}} (p^\nu e_\nu) = e_\nu [\nabla_\alpha, \nabla_\beta] p^\nu = e_\nu (R_{\alpha\beta\mu}{}^\nu p^\mu - \Gamma_{\alpha\beta}^\sigma \nabla_\sigma p^\nu)$$

② Autoparallels are the Papapetrou Eqs

$$0 = \frac{d}{d\kappa} (e_\mu p^\mu) = e_\mu \left[\dot{\chi}^\sigma \nabla_\sigma + \frac{1}{2} \dot{\alpha}^{\alpha\beta} [\nabla_\alpha, \nabla_\beta] + \frac{d}{d\kappa} \right] p^\mu$$

$$0 = \frac{d}{d\kappa} (e_{\mu\nu} S^{\mu\nu}) = e_{\mu\nu} \left[\dot{\chi}^\sigma \nabla_\sigma + \frac{1}{2} \dot{\alpha}^{\alpha\beta} [\nabla_\alpha, \nabla_\beta] + \frac{d}{d\kappa} \right] S^{\mu\nu}$$

③ TERMS in Red are Non Geodesic, Violate EEP

$$\dot{p}^\mu + p^\nu \dot{\chi}^\beta \Gamma_{\beta\nu}^\mu + \frac{1}{2} \dot{\chi}^\sigma S^{\alpha\beta} R_{\alpha\beta\sigma}{}^\mu = 0$$

$$\dot{S}^{\alpha\beta} + \dot{\chi}^\nu (S^{\delta\beta} \Gamma_{\nu\delta}^\alpha + S^{\alpha\delta} \Gamma_{\nu\delta}^\beta)$$

$$+ \frac{1}{2m} S^{\mu\nu} (S^{\delta\beta} R_{\mu\nu\delta}{}^\alpha + S^{\alpha\delta} R_{\mu\nu\delta}{}^\beta) = 0$$

8 M.C Curved Space Lagrangian

Spinning Particles follow paths which are NOT 4D geodesics but ARE autoparallels in 'polygeometric space'

① Lagrangian - Just add metric!

$$\mathcal{L} = m_0 \sqrt{\dot{x}^\mu \dot{x}^\nu g_{\mu\nu} - \frac{1}{2\lambda^2} \dot{a}^{\mu\alpha} \dot{a}^{\nu\beta} g_{\mu\nu} g_{\alpha\beta}}$$

② Dirac Canonical Form (Better to quantize) vary p independent of x ...

$$\mathcal{L} = p_\mu \dot{x}^\mu - \frac{1}{2\lambda^2} S_{\mu\nu} \dot{a}^{\mu\nu} - \frac{1}{2m_0} \left(p_\mu p^\mu g^{\alpha\beta} - \frac{S_{\mu\alpha} S_{\nu\beta} g^{\mu\nu} g^{\alpha\beta}}{2\lambda^2} \right)$$

③ 3rd Form Vary δm_0 as Lagrange Multiplier

$$\mathcal{L} = m_0 \left(\dot{x}_\mu \dot{x}^\mu - \frac{\dot{a}_{\mu\nu} \dot{a}^{\mu\nu}}{2\lambda^2} \right) + \frac{1}{m_0} \left(p_\mu p^\mu - \frac{S_{\mu\nu} S^{\mu\nu}}{2\lambda^2} \right)$$

④ Possible Extensions: add mass as canonical to scalar (+ weyl connection)

$$\mathcal{L} = \sqrt{-\dot{\phi}^2 + \dot{x}^2 - \frac{1}{2\lambda^2} \dot{a}^2}$$

$$= -m_0 \dot{\phi} + p_\mu \dot{x}^\mu - \frac{1}{2\lambda^2} S_{\mu\nu} \dot{a}^{\mu\nu} - \lambda \left(m_0^2 - p^2 + \frac{S^2}{2\lambda^2} \right)$$

↑
multiplier

IV A Anholonomic Mechanics

① Variation of Lagrangian

$$\delta \mathcal{L} = \left(\frac{\delta \mathcal{L}}{\delta x^\alpha} \delta x^\alpha + \frac{\delta \mathcal{L}}{\delta \dot{x}^\alpha} \delta \dot{x}^\alpha \right) + \frac{1}{2} \left(\frac{\delta \mathcal{L}}{\delta a^{\alpha\beta}} \delta a^{\alpha\beta} + \frac{\delta \mathcal{L}}{\delta \dot{a}^{\alpha\beta}} \delta \dot{a}^{\alpha\beta} \right)$$

$$\rightarrow p_\alpha \delta \dot{x}^\alpha = \frac{d}{d\tau} \left(p_\alpha \delta x^\alpha \right) - \underbrace{\dot{p}_\alpha \delta x^\alpha}_{\text{Standard}} + p_\alpha \left(\delta \dot{x}^\alpha - \frac{d}{d\tau} \delta x^\alpha \right)_{\text{Does Not vanish!}}$$

② Variation of Velocity Component is **Not** the velocity of Variation Component

Shortcuts: $\delta(d x^\mu e_\mu) = d(\delta x^\mu e_\mu)$

$$\rightarrow e_\mu \left(\delta \dot{x}^\alpha - \frac{d}{d\tau} \delta x^\alpha \right) = \left(\delta x^\alpha \dot{e}_\alpha - \dot{x}^\alpha \delta e_\alpha \right)$$

$$= e_\sigma \left[\delta x^\alpha \dot{x}^\beta \tau_{\alpha\beta}^\sigma + \frac{1}{2} \left(\delta x^\alpha \dot{a}^{\mu\nu} - \dot{x}^\alpha \delta a^{\mu\nu} \right) R'_{\mu\nu\alpha}{}^\sigma \right]$$

$$\rightarrow \left(\delta \dot{a}^{\alpha\beta} - \frac{d}{d\tau} \delta a^{\alpha\beta} \right) = 2 \left(\delta a^{\sigma\rho} \dot{x}^\alpha - \dot{a}^{\sigma\rho} \delta x^\alpha \right) \Gamma_{\lambda\sigma}^\alpha + \left(\delta a^{\sigma\rho} \dot{a}^{\mu\nu} - \dot{a}^{\sigma\rho} \delta a^{\mu\nu} \right) R'_{\mu\nu\sigma}{}^\alpha$$

Note $R'_{\mu\nu\sigma}{}^\alpha \equiv R_{\mu\nu\sigma}{}^\alpha - \tau_{\mu\nu}{}^\lambda \Gamma_{\lambda\sigma}^\alpha$

③ With this, get Papapetrou eqns!

$$\dot{p}_\alpha - \frac{\delta \mathcal{L}}{\delta x^\alpha} = p_\sigma \tau_{\alpha\beta}^\sigma \dot{x}^\beta + p_\sigma R'_{\mu\nu\alpha}{}^\sigma \dot{a}^{\mu\nu} + S_{\sigma\beta} \tau_{\alpha\mu}^\sigma \dot{a}^{\mu\beta}$$

IV.B Hamiltonian Canonical Mechanics Useful for Quantizing!

① Legendre Transformation $\mathcal{H} = p_\alpha \dot{x}^\alpha - \frac{1}{2\lambda^2} S_{\alpha\beta} \dot{a}^{\alpha\beta} - \mathcal{L}$

$$\begin{aligned} \delta \mathcal{H} &= \left(\frac{\partial \mathcal{H}}{\partial p_\alpha} \delta p_\alpha + \frac{\partial \mathcal{H}}{\partial x^\alpha} \delta x^\alpha \right) + \frac{1}{2} \left(\frac{\partial \mathcal{H}}{\partial S_{\alpha\beta}} \delta S_{\alpha\beta} + \frac{\partial \mathcal{H}}{\partial a^{\alpha\beta}} \delta a^{\alpha\beta} \right) \\ &= (\dot{x}^\alpha \delta p_\alpha - \frac{\partial \mathcal{L}}{\partial x^\alpha} \delta x^\alpha) - \frac{1}{2\lambda^2} (\dot{a}^{\alpha\beta} \delta S_{\alpha\beta} + \frac{\partial \mathcal{L}}{\partial a^{\alpha\beta}} \delta a^{\alpha\beta}) \end{aligned}$$

② Non commutativity of $[\delta, d] \neq 0$ destroys symmetry of Hamilton's Equations

ok $\dot{x}^\mu = \frac{\delta \mathcal{H}}{\delta p_\mu} \quad \dot{a} = -\lambda^2 \frac{\delta \mathcal{H}}{\delta S_{\mu\nu}}$

BUT

$$\begin{aligned} \frac{\delta \mathcal{H}}{\delta x^\lambda} &= -\dot{p}_\lambda + p_\sigma T_{\lambda\rho}^\sigma \frac{\delta \mathcal{H}}{\delta p_\rho} + p_\sigma R'_{\mu\nu\lambda}{}^\sigma \frac{\delta \mathcal{H}}{\delta S_{\mu\nu}} \\ &\quad - S_{\sigma\beta} T_{\lambda\mu}^\sigma \frac{\delta \mathcal{H}}{\delta S_{\mu\nu}} \end{aligned}$$

③ Poisson Brackets is $\{F, G\} \rightarrow \frac{1}{i\hbar} [\hat{F}, \hat{G}]$ valid?

$$\dot{F} = \left(\frac{\partial F}{\partial x} \dot{x} + \frac{\partial F}{\partial p} \dot{p} \right) + \frac{1}{2} \left(\frac{\partial F}{\partial a} \dot{a} + \frac{\partial F}{\partial S} \dot{S} \right)$$

$$= \{F, \mathcal{H}\} + \text{Torsion} + \text{Curvature Terms}$$

$$\rightarrow \left(\frac{\partial F}{\partial x^\alpha} \frac{\partial \mathcal{H}}{\partial p_\mu} - \frac{\partial F}{\partial p_\mu} \frac{\partial \mathcal{H}}{\partial x^\alpha} \right) - \frac{\lambda^2}{2} \left(\frac{\partial F}{\partial a^{\mu\nu}} \frac{\partial \mathcal{H}}{\partial S_{\mu\nu}} - \frac{\partial F}{\partial S_{\mu\nu}} \frac{\partial \mathcal{H}}{\partial a^{\mu\nu}} \right)$$

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